

Enhanced reverse saturable absorption and optical limiting in heavy-atom-substituted phthalocyanines

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The reverse saturable absorption and the optical-limiting response of metal phthalocyanines can be enhanced by use of the heavy-atom effect. Phthalocyanines containing heavy-metal atoms, such as In, Sn, and Pb, show a nearly factor-of-2 enhancement in the ratio of effective excited-state to ground-state absorption cross sections compared with those containing lighter atoms, such as Al and Si. In an $f/8$ optical geometry, homogeneous solutions of heavy-metal phthalocyanines, at 30% linear transmission, limit 8-ns 532-nm laser pulses to $\leq 3 \mu\text{J}$ (the energy for 50% probability of eye damage) for incident energies as high as 800 μJ .

Materials that exhibit reverse saturable absorption (RSA) are currently of interest for use in optical-limiting devices for protection of sensors and eyes from energetic light pulses. RSA can occur when states with an absorption cross section (σ_e) in excess of the ground-state cross section (σ_g) are produced. RSA in the visible spectrum that is due to electronic excitation was reported for several classes of organic chromophores.¹⁻⁷ The effectiveness of RSA molecules for optical limiting is determined mainly by the ratio of cross sections (σ_e/σ_g), which is function of wavelength, and the populations of the states that evolve in time during the pulse. There is considerable interest in approaches to enhancing the performance of RSA chromophores.

RSA and optical limiting in metallophthalocyanines (MPc's) such as chloroaluminum phthalocyanine (CAP) and bis[tri-(*n*-hexyl)siloxy] silicon naphthalocyanine (SiNc) have been investigated^{4,7} on picosecond and nanosecond time scales at 532 nm. Picosecond experiments⁷ showed that values for σ_s/σ_g , where s refers to the excited-singlet state, were 10.5 and 14 for CAP and SiNc, respectively. Nanosecond measurements on CAP (Ref. 8) and SiNc (Ref. 9) indicate that σ_t/σ_g values (t refers to the triplet state) are ~ 20 to 50. However, the triplet populations achievable during nanosecond pulses are limited by the small intersystem crossing rate, k_{isc} , for these molecules.^{9,10} These observations suggested to us the use of the heavy-atom effect as an approach to enhancing the optical-limiting performance of phthalocyanines. When the atomic number of the central metal atom is increased, k_{isc} for the π -electronic states can be increased,¹⁰ permitting more efficient population of the triplet state.

In this Letter we examine the RSA performance of MPc's containing metals (and metalloids) from groups IIIA (Al, Ga, In) and IVA (Si, Ge, Sn, and Pb). For Al, Ga, and In phthalocyanines (Pc's) the molecular structures were of the form [tri-(*n*-hexyl)siloxy] MPc,

whereas Si, Ge, and Sn Pc's were of the form bis[tri-(*n*-hexyl)siloxy] MPc and Pb Pc was of the form Pb tetra(*t*-butyl)Pc. These molecules were synthesized according to literature methods.¹¹ The electronic absorption spectrum of Sn Pc is shown in Fig. 1, which reveals a strong *Q*-band absorption at 678 nm and a region of weak absorption from ~ 425 to 600 nm. Also shown is the transient triplet-triplet absorption spectrum, which shows a maximum at 510 nm and a bandwidth of ~ 130 nm. The ground-state and triplet-triplet spectra for the other molecules are quite similar, with a slight red shift of the *Q* band as the metal becomes heavier.

The values of the triplet quantum yield (Φ_t) and the first excited-singlet lifetime (τ_s) are given in Table 1 for the group IIIA and group IVA MPc's. Indeed, the τ_s values decrease and the Φ_t values increase as

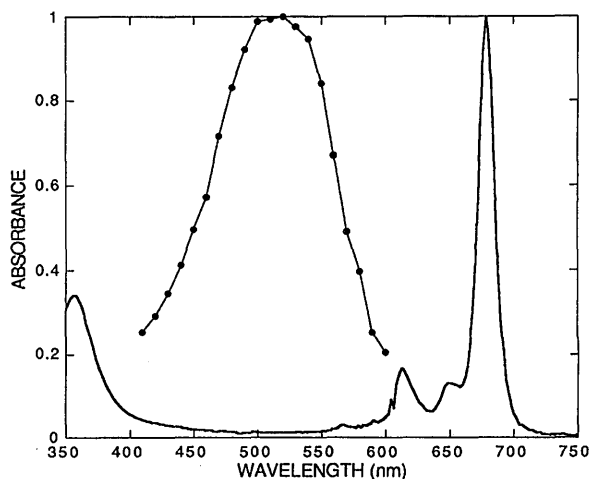


Fig. 1. Electronic absorption spectrum of bis[tri-(*n*-hexyl)siloxy]SnPc in toluene solution at 7.1×10^{-6} M. Also shown (curve with circles) is the transient absorption spectrum (arbitrary units) of Sn Pc in toluene obtained 100 ns after excitation at 355 nm.

Table 1. Photophysical Properties for Group IIIA and Group IVA Metal Pc's

Metal	τ_s (ns) ^a	Φ_t ^b	f_t (8 ns)
Al	6.9	0.35	0.18
Ga	3.3	0.51	0.42
In	0.3	0.88	0.95
Si	4.5 ^c	0.35 ^d	0.25
Ge	4.2 ^c	0.37 ^d	0.28
Sn	2.0 ^c	0.62	0.63
Pb	0.35	0.92	0.95

^aFirst excited-state lifetimes determined by picosecond pump-probe measurements at 700 nm; uncertainty is $\pm 10\%$.

^bTriplet quantum yields determined from ground-state bleaching recovery at 700 nm; uncertainty is ± 0.10 .

^cValues determined by use of time-correlated photon counting.

^dValues estimated from fluorescence quantum yields.

the metal becomes heavier. Thus the triplet population achievable during a *Q*-switched laser pulse (e.g., 8-ns) is significantly enhanced for the molecules with large Φ_t values. To the extent that the different molecules have similar σ_g , σ_s , and σ_t values, one would expect trends of increased nonlinear absorption for the heavier-metal Pc's, for pulse widths between ~ 0.5 and 25 ns.

Nanosecond nonlinear transmission measurements were performed with a frequency-doubled *Q*-switched Nd:YAG laser operating in a near-Gaussian transverse mode with an 8-ns (FWHM) multimode pulse envelope. The laser pulses were focused at the center of 1-cm path-length solution cells by a 15-cm focal-length, 2.5-cm-diameter best-form lens (*f*/40 geometry) giving a waist spot size of $\sim 13 \mu\text{m}$ (half-width at $1/e^2$ maximum intensity) and a cell-entrance spot size of $\sim 63 \mu\text{m}$. The transmitted energy was collected by a 2.5-cm-diameter, 15-cm focal-length lens located so as to relay the collected energy to a nominal spot size of ~ 2.0 mm on a photodiode of 1-cm diameter. This geometry afforded essentially total integrated energy detection, as was verified by the following experiment. Measurements with a 50%-transmitting aperture stop in front of the detector performed on SiNc in a series of solvents with varying thermal refractive-index coefficients (i.e., nitrobenzene to carbon disulfide) showed different transmission responses; whereas, without the aperture, the responses were the same, thus demonstrating the insensitivity of the unapertured detection to refractive spot-size changes for pulse energies as high as the highest values used.

Nonlinear transmission data obtained with 532-nm, 8-ns laser pulses for the two series of MPC's are shown in Fig. 2. At incident fluences below $\sim 2 \text{ mJ/cm}^2$ the transmittance coincides with that measured by use of a spectrophotometer. The nanosecond pulse data for the two series of molecules show the trend of increased nonlinear absorption as the metal becomes heavier, consistent with an enhanced triplet contribution. Measurements on the group IVA MPC's were also performed with 70-ps pulses, which are short compared with the time scales for triplet formation in these molecules. The Si, Ge, and Sn Pc's showed the opposite trend of

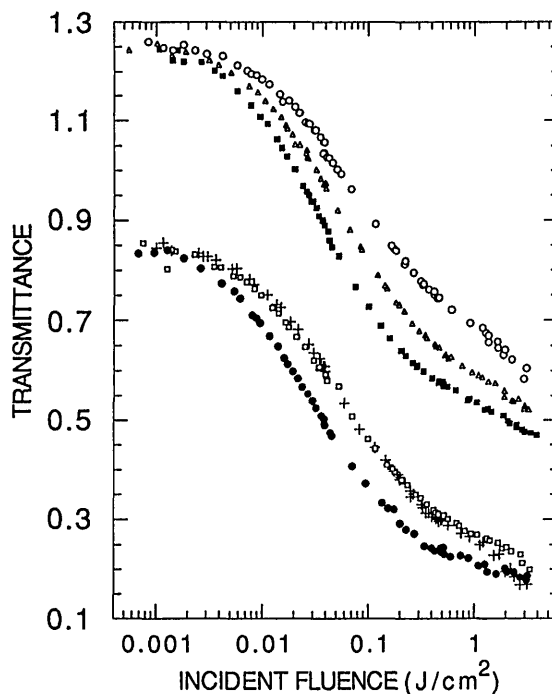


Fig. 2. Nonlinear transmittance of group IIIA (upper curves: Al, open circles; Ga, open triangles; In, filled squares) and group IVA (lower curves: Si, plus sign; Ge, open squares; Sn, filled circles) metal-substituted Pc's in toluene solution measured with 8-ns, 532-nm laser pulses and an *f*/40 optical geometry. Solutions had 84% linear transmission at 532 nm for a 1-cm path length, which corresponds to a concentration of $\sim 1.4 \times 10^{-4}$ M. The upper curves have been displaced vertically by 0.4 for presentation.

the nonlinear transmittance compared with that for the nanosecond pulses. Since the σ_g values are all $\sim 2.3 \times 10^{-18} \text{ cm}^2$ for these molecules the picosecond results show that the relative strength of the excited-singlet absorption is opposite the apparent absorption observed on the nanosecond time scale. These results are consistent with enhanced triplet populations for the heavier-metal molecules during the longer pulses.

The nonlinear transmission functions for reverse saturable absorbers are expected to range from $T_{\text{lin}} = \exp(-\sigma_g N_0 L)$ for low incident fluences to a saturated transmission $T_{\text{sat}} = \exp(-\sigma_e N_0 L)$ for high degrees of excitation. When the pulse width is short compared with the excited-state (say, the triplet) lifetime, we can define a saturation fluence $F_{\text{sat}} = h\nu/\sigma_g \Phi_t$. At such high fluences the ground-state population is largely depleted, and the excited-state population is distributed between the first excited-singlet and lowest-triplet states. A figure of merit for RSA molecules involving excited-singlet and triplet states can be defined as

$$\sigma_{\text{eff}}/\sigma_g = \ln T_{\text{sat}}/\ln T_{\text{lin}} \approx (\sigma_s f_s + \sigma_t f_t)/\sigma_g,$$

where σ_{eff} is an effective excited-state cross section and f_s and f_t are average fractional populations of the excited-singlet and triplet states during the pulse.

From Fig. 2 we estimate that $T_{\text{sat}} \sim 0.07$ for In Pc. For this molecule, f_t (Table 1) is ~ 0.95 , so we

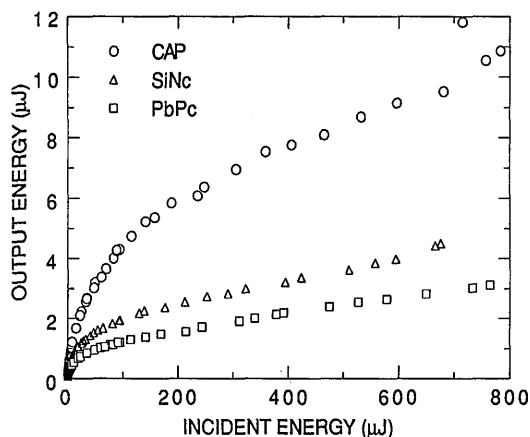


Fig. 3. The $f/8$ optical-limiting response of Pb Pc and SiNc in toluene and of CAP in methanol at 532 nm. The input lens had a 2.5-cm diameter and 4-cm focal length; the laser beam diameter was 4.8 mm (full width at $1/e^2$ maximum intensity) at the lens; the calculated beam waist was $10\text{ }\mu\text{m}$ (half-width at $1/e^2$ maximum intensity); and the focus was close to the exit window of the 1-cm cell. The transmitted energy was collected at $\sim f/2$ and relayed to a 1.5-mm spot size on a 1-cm-diameter detector. Solutions had a linear transmittance of 0.3 at 532 nm; concentrations were $7 \times 10^{-4}\text{ M}$ for SiNc, $9 \times 10^{-4}\text{ M}$ for CAP, and $1.3 \times 10^{-3}\text{ M}$ for Pb Pc.

can ignore f_s . Using a typical σ_t value of $50 \times 10^{-18}\text{ cm}^2$ for MPc's, we calculate $T_{\text{sat}} = 0.02$. This value is smaller than that observed, yet the fluence exiting the sample was approximately equal to F_{sat} (0.24 J/cm^2 for In Pc), suggesting that uniform on-axis saturation was being approached. In any case we can use the high-fluence transmittance to place a lower bound on $\sigma_{\text{eff}}/\sigma_g$ for the various molecules. The cross-section ratios range from 10 to 16 for the group IIIA MPc's and from 10 to 18 for the group IVA MPc's as the metal becomes heavier. In comparison, the fullerene C_{60} was reported¹² to have a σ_t/σ_g value of ~ 3 at 532 nm and a Φ_t value of ≥ 0.96 .

We have investigated the performance of the heavy-atom Pc's in a simple optical-limiting geometry. Figure 3 shows the optical-limiting response in an $f/8$ geometry of Pb Pc compared with those of CAP and SiNc. To our knowledge, SiNc had been the strongest known absorptive optical-limiting dye for nanosecond 532-nm pulses. The strong signal output of Pb Pc is lower than that of CAP and SiNc by factors of 4 and 1.5, respectively. More importantly, the maximum allowable input energy for $3\text{-}\mu\text{J}$ output energy, i.e., the energy corresponding to 50% probability for retinal damage to eyes, is higher than that for CAP and SiNc by factors of 16 and 2, respectively.

These results demonstrate the enhanced capability (i.e., nonlinear suppression of nanosecond 532-nm pulses by a factor of 81 for a limiter with a linear transmission of 30% in an $f/8$ optical system) of the heavy-atom Pc's for optical-limiting applications. Recently strong optical limiting was also reported for a Pb(tetrakis-cumylphenoxy)Pc derivative.¹³ It should be emphasized that the optical-limiting geometries that have been used are far from optimal. Recent device modeling¹⁴ suggests that heavy-atom Pc materials in optimized optical-limiter designs that ensure uniform saturation could produce attenuations of 10^4 for devices with 70% linear transmission.

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